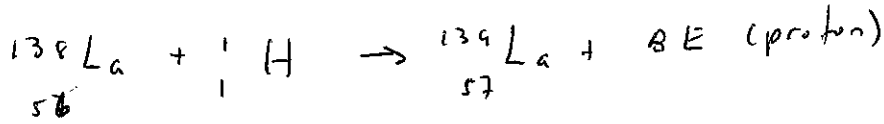
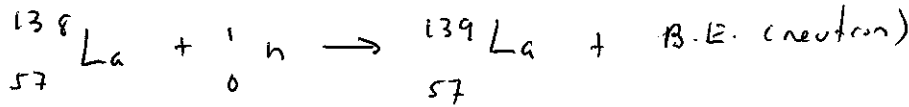


Chapter 1

1 (i) Binding energy of a proton and neutron in  $^{139}\text{La}$ :



$$\text{BE (neutron)} = \left[ {}^{138}_{57}\Delta + \Delta n - {}^{139}_{57}\Delta \right] = \left[ -86524.681 + 8071.3171 - (-87231.371) \right] \\ = 8778.0071 \text{ keV}$$

$$\text{BE (proton)} = \left[ {}^{138}_{56}\Delta + \Delta H - {}^{139}_{57}\Delta \right] = \left[ -88261.631 + 7288.9705 - (-87231.371) \right] \\ = 6257.7105 \text{ keV}$$

$\therefore$  The binding energy of the neutron =  $\frac{8.778(0071) \text{ MeV}}{6.257(7105)}$  and the binding energy of the proton } This difference is related to structure  $\rightarrow$  odd # of protons and even # of neutrons (i.e., 82 magic #)

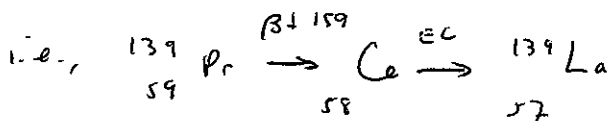
(ii) The mean binding energy per nucleon is:

$$\text{For } {}^{139}_{57}\text{La} = \frac{[Z\Delta_H + (A-Z)\Delta_n - \frac{A}{Z}\Delta]}{A} \frac{\text{keV}}{\text{nucleon}} \\ = \frac{[57 \times 7288.9705 + (139-57) 8071.3171 - (-87231.371)]}{139} = 8378.0625 \frac{\text{keV}}{\text{nucleon}}$$

$\therefore$  The mean binding energy is  $\frac{8.378063 \text{ MeV}}{\text{nucleon}}$  (same number as Appendix B)

This value is slightly less than that given for the neutron but greater than that for the proton since the neutron is paired in La-139 whereas the proton is the odd one.

(iii)  ${}^{139}_{59}\text{Pr}$  and  ${}^{139}_{58}\text{Ce}$  have  $\beta^+$  and EC transition energies of 2.11 and 0.28 MeV, respectively.



$$\text{Now } E(\beta^+) = (\text{mass } {}^{139}\text{Pr} - \text{mass } {}^{139}\text{Ce}) \times 931.494 - 2 \times 0.511 \text{ MeV and} \\ E(\text{EC}) = (\text{mass } {}^{139}\text{Ce} - \text{mass } {}^{139}\text{La}) \times 931.494 \text{ MeV (using CODATA data).}$$

$$\therefore 2.11 = E(\beta^+) = (\text{mass } {}^{139}\text{Pr} - \text{mass } {}^{139}\text{Ce}) \times 931.494 - 2 \times 0.511 \text{ MeV} \quad (1)$$

$$0.28 = E(\text{EC}) = (\text{mass } {}^{139}\text{Ce} - \text{mass } {}^{139}\text{La}) \times 931.494 \text{ MeV} \quad (2)$$

Mass excess of  ${}^{139}_{57}\text{La}$  is  ${}^{139}_{57}\Delta = -87231.371 \text{ keV}$  or in mass units:

$${}^{139}\Delta = \frac{-87.231371 \text{ MeV}}{931.494 \frac{\text{MeV}}{\text{amu}}} = -93646.73 \mu\text{u}$$

1-3

$$\therefore {}^{139}\text{Mass La} = 139 - 93646.73 \times 10^{-6} = \underline{138.90635 \text{ u}}$$

$$\text{Using Eq. (2): } \text{Mass } {}^{139}\text{Ce} = \frac{.28}{931.494} + 138.90635 = \underline{138.90665 \text{ u}}$$

$$\text{and Eq. (1): } \text{Mass } {}^{139}\text{Pr} = \frac{2.11 + 2(0.511)}{931.494} + 138.90665 = \underline{138.91001 \text{ u}}$$

(iv) Consider the binding energy equation:

$$B_{A,Z} = [Z m_H + (A-Z) m_n - M_{A,Z}] 931.494 \text{ MeV or in terms of mass defect}$$

$$\delta M_{A,Z} = \frac{B_{A,Z}}{931.494} = [Z m_H + (A-Z) m_n - M_{A,Z}]$$

Given the definition of mass excess  $M_{A,Z} = \Delta A + A$

$$\begin{aligned} \therefore \delta M_{A,Z} &= [Z (\Delta H + 1) + (A-Z) (\Delta n + 1) - (\Delta A + A)] \\ &= [Z \Delta H + Z + A \Delta n + A - Z \Delta n - Z - \Delta A - A] \end{aligned}$$

$$\therefore \delta M_{A,Z} = [Z \Delta H + (A-Z) \Delta n - \Delta A]$$

(v) For  ${}^{139}_{59}\text{Pr}$ , A is odd  $\Rightarrow$  spin + parity obtained from the 59<sup>th</sup> proton

in the  $2d_{5/2}$  level  $\therefore$  for d,  $\ell=2 \therefore +$  parity and spin  $\frac{5}{2}$

Similarly for  ${}^{139}_{58}\text{Ce}$ , the spin and parity are obtained from the

81<sup>st</sup> neutron in the  $1h_{11/2}$  level  $\Rightarrow$  for h,  $\ell=5 \therefore -$  parity and spin  $\frac{11}{2}$

(actually  $\frac{3}{2}$ )

For  ${}^{139}_{57}\text{La}$ , the spin and parity are obtained from the 57<sup>th</sup> proton in the

$1g_{7/2}$  level  $\Rightarrow$  for g,  $\ell=4 \therefore +$  parity and spin  $\frac{7}{2}$ .

However, the  $\Delta J$  for the positron transition from the ground states for

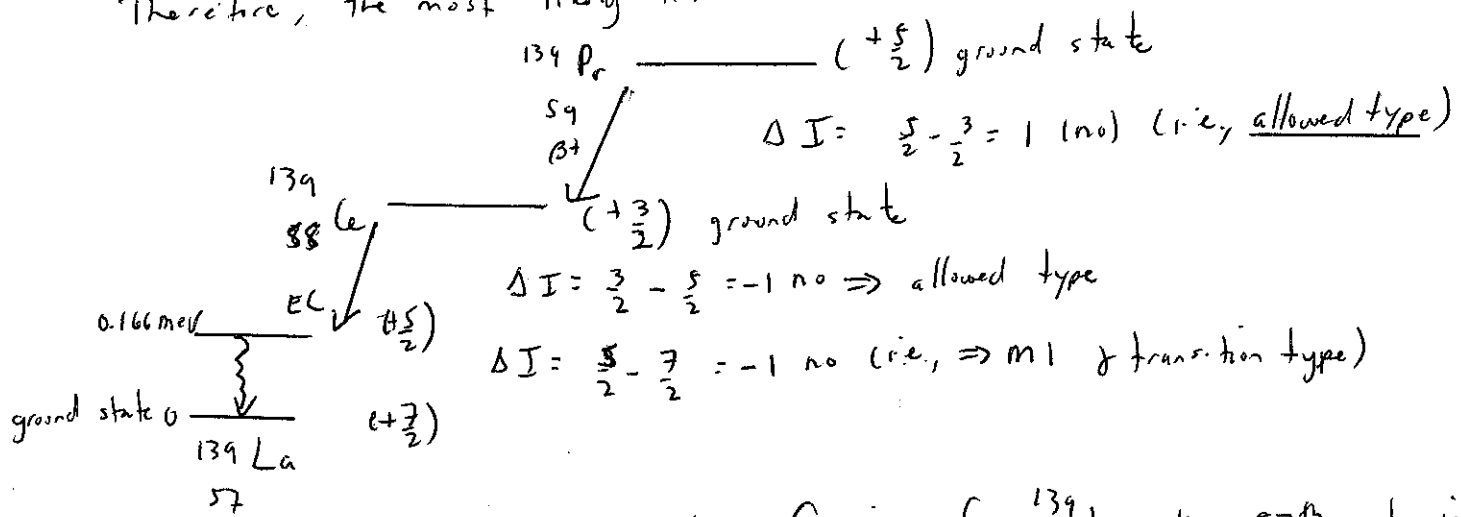
$${}^{139}_{59}\text{Pr} \left( +\frac{5}{2} \right) \xrightarrow{\beta^+} {}^{139}_{58}\text{Ce} \left( -\frac{11}{2} \right) \text{ is } \Delta J = \frac{5}{2} - \frac{11}{2} = -3 \text{ yes (there is a parity change)}$$

Hence the beta selection rules suggest a 3<sup>rd</sup> forbidden type.

However, since we are near a magic number for  ${}^{139}\text{Ce}$  (81<sup>st</sup> neutron) and

the fact that the spin prediction suggest a 3rd forbidden type, it is reasonable that a lower energy state may exist by promoting a nucleon from a level of lower angular momentum into the higher angular momentum level so that a pair can be formed. Hence, for  $^{139}\text{Ce}$ , it appears that a neutron removed from a pair in the  $2d_{3/2}$  level to form a pair with the  $1h_{1/2}$  neutron; the odd particle is therefore a  $2d_{3/2}$  neutron (implying a  $+3/2$  parity and spin ground state. (Alternatively, a neutron could have been promoted from the  $3s_{1/2}$  state for  $^{139}\text{Ce}$ , however this would have implied a greater angular momentum change for the positron transition ( $\Delta I = \frac{5}{2} - \frac{1}{2} = 2$  no  $\Rightarrow$  2nd forbidden type), as well as for the EC transition.

Therefore, the most likely transitions are:



Therefore, with reference to the above figure, for  $^{139}\text{La}$  the 57th neutron is in the  $1g_{7/2}$  level (ground state  $+7/2$ ) and in the energy level diagram the next state is  $2d_{5/2}$  (i.e., 1st excited state  $+5/2$ ). Hence, the  $\Delta I$  change for the EC transition is less between the ground state of  $^{139}\text{Ce}$  and the 1st excited state of  $^{139}\text{La}$ , than between the two ground states, implying a greater decay to the excited state. Note that the  $\Delta I$  change for this E.C. transition is  $\Delta I = \frac{3}{2} - \frac{5}{2} = -1$  no, which according to  $\beta$  decay systematics is an allowed type hence this transition is likely by E.C. decay since the theory of EC is similar to  $\beta$  decay (i.e., it only involves the further wave function of the atomic electrons). Further,  $^{139}\text{Ce}$  undergoes an EC decay rather than  $\beta^+$  decay because the mass which equals 0.28 MeV is less than the threshold 1.022 MeV for  $\beta^+$  decay.

(vi) Using the Binding energy per nucleon (Appendix B) for  $^{139}_{59}\text{Pr}$ ,  $^{139}_{58}\text{Ce}$ ,  $^{139}_{57}\text{La}$ , the corresponding binding energies  $B$  are:

$$B_{\text{Pr}} = 8349.482 \text{ (keV/nucleon)} \times 139 \text{ nucleons} = 1160.578 \text{ MeV}$$

$$B_{\text{Ce}} = 8370.428 \text{ (keV/nucleon)} \times 139 \text{ nucleons} = 1163.489 \text{ MeV}$$

$$B_{\text{La}} = 8378.063 \text{ (keV/nucleon)} \times 139 \text{ nucleons} = 1164.551 \text{ MeV}$$

Given the Weizsäcker formula for odd  $A$  nucleides:

$$B(A, Z) = 14.1A - 13.1A^{2/3} - 0.58Z(Z-1)A^{-1/3} - 18(N-Z)^2A^{-1}$$

Hence, the binding energies are:

	$^{139}_{59}\text{Pr}$	$^{139}_{58}\text{Ce}$	$^{139}_{57}\text{La}$
- $B$ (mass data) (MeV)	-1160.578	-1163.489	-1164.551
- $B(A, Z)$ (Weizsäcker) (MeV)	-1168.12	-1169.72	-1170.05
$Z$	59	58	57

The discrepancy in binding energies is due to the fact that the constants in the Weizsäcker formula are derived over a large range of  $A$  values. Thus, extrapolating the binding energies by fitting the constants to the specific mass data where:

$$B = aZ^2 + bZ + c$$

for the matrix problem  $A\vec{x} = \vec{b}$

where  $A = \begin{bmatrix} 59^2 & 59 & 1 \\ 58^2 & 58 & 1 \\ 57^2 & 57 & 1 \end{bmatrix}$ ,  $\vec{x} = \begin{bmatrix} a \\ b \\ c \end{bmatrix}$ ,  $\vec{b} = [-1160.578, -1163.489, -1164.551]$

can be solved by Maple (see attached) yielding the solution

$$a = 0.9244995555, \quad b = -105.2454489, \quad c = 1831.310530$$

See attached plot of  $-B$  (meV) versus  $Z$

```
[ > with(linalg):  
[ > b:= vector([-1160.578,-1163.489,-1164.551]);  
      b := [-1160.578,-1163.489,-1164.551]  
[ > A:= matrix(3,3,[59^2,59,1,58^2,58,1,57^2,57,1]);  
      A :=  $\begin{bmatrix} 3481 & 59 & 1 \\ 3364 & 58 & 1 \\ 3249 & 57 & 1 \end{bmatrix}$   
[ > x:= linsolve(A,b);  
      x := [0.9244995555,-105.2554489,1831.310530]  
[ > a:=.9244995555; b:=-105.2554489;c:=1831.310530;  
      a := 0.9244995555  
      b := -105.2554489  
      c := 1831.310530  
[ > plot((a*Z^2+b*Z+c), Z = 53..59);#(-BE (MeV) vs Z)
```

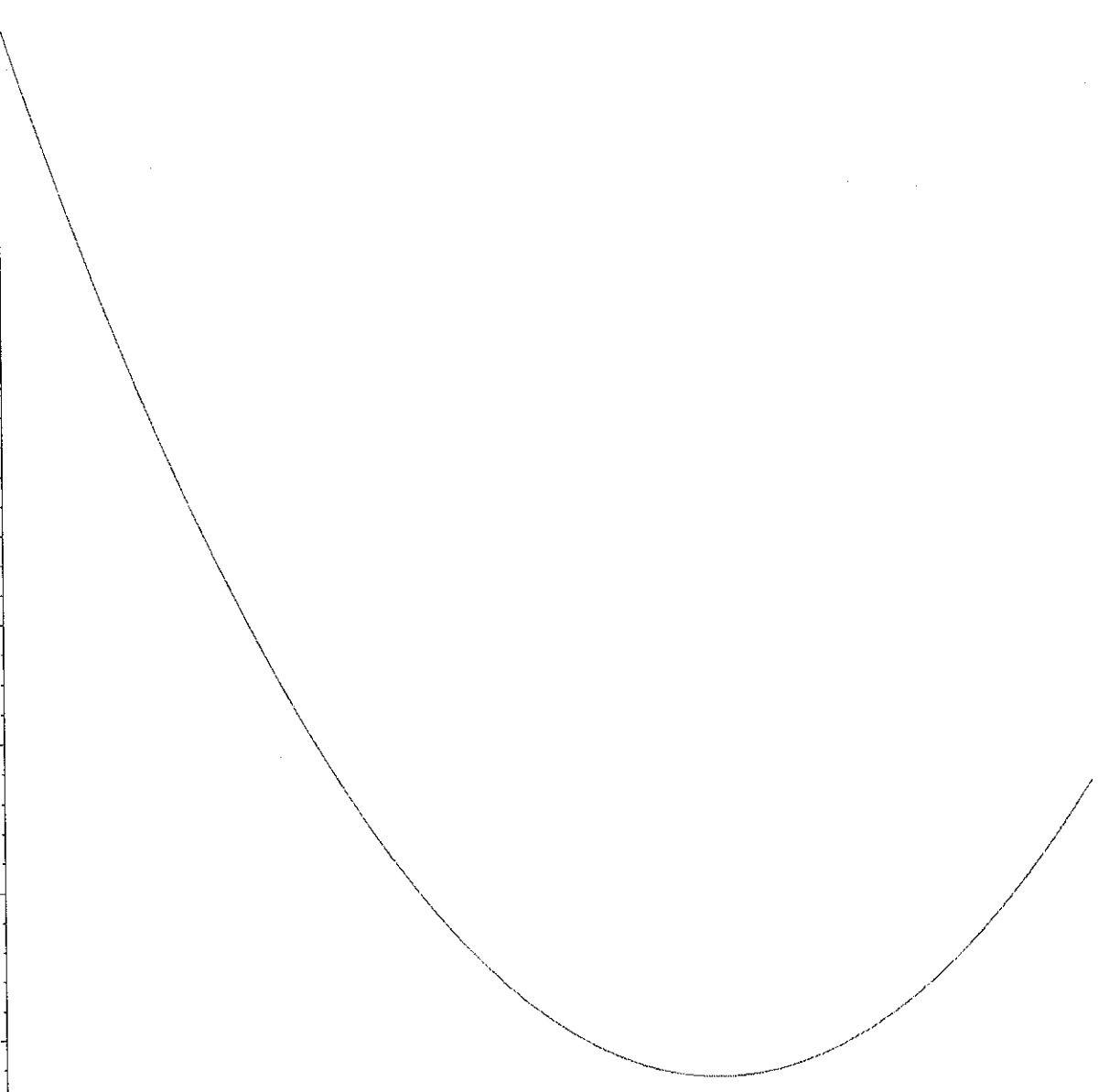
-B (mev)

-1152  
-1154  
-1156  
-1158  
-1160  
-1162  
-1164

53 54 55 56 57 58 59

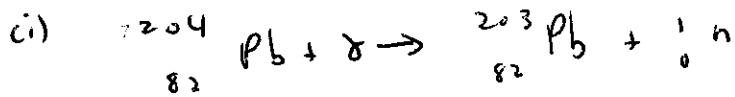
Z

[ >  
[ >



2. The threshold energy  $T_m$  for a particle of mass  $m$  striking a nucleus of mass  $M$  with a reaction energy  $Q$  is: 1-6

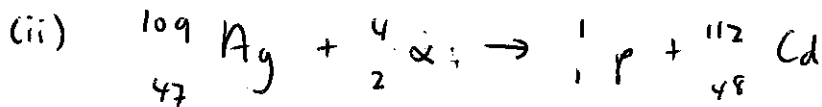
$$T_m = -Q \frac{(M+m)}{M}$$



$\therefore Q = (\Delta^{204}\text{Pb} - \Delta^{203}\text{Pb} - \Delta\text{n})$  where  $\Delta$  is the mass excess in keV:

$$Q = (-25109.735 - (-24786.57) - 8071.3171) = -8394.482 \text{ keV}$$

This reaction is endothermic and the threshold energy necessary to make this reaction go is (since  $m=0$  for a  $\gamma$ -ray)  $\Rightarrow T_m = -Q = \underline{8.394 \text{ MeV}}$



$\therefore Q = (\Delta^{109}\text{Ag} + \Delta^4\text{He} - \Delta^{112}\text{Cd} - \Delta^1\text{H}) \text{ keV}$  (where 2 electron masses were added to each side of the equation)

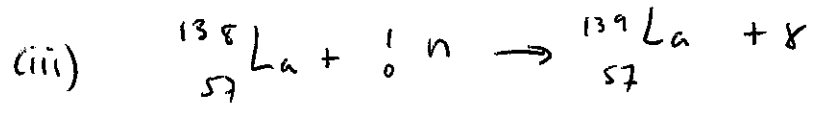
$$\therefore Q = (-88722.669 + 2424.91565 - (-90580.518) - 288.9705) \text{ keV} = -3006.206 \text{ keV} = \underline{-3.006 \text{ MeV}}$$

This reaction is also endothermic. Therefore, the threshold energy is:

$$T_m = -Q \frac{(M+m)}{M} = +3.006206 \frac{(108.904752 + 4.001506)}{108.904752} = \underline{3.117 \text{ MeV}}$$

Here  $m = 4.00260325 \text{ u} - 2 \times \frac{511006 \text{ MeV}}{931.494 \text{ MeV/u}} = 4.001506 \text{ u}$

The atomic masses are taken from Appendix B and  $m$  is the mass of the He nucleus ( $\alpha$  particles).



$$Q = [\Delta({}^{138}\text{La}) + \Delta n - \Delta({}^{139}\text{La})] \text{ keV}$$

$$\therefore Q = [-86524.681 + 8071.3171 - (-87231.371)] \text{ keV} = 8778.0071$$

$$= \underline{8.778 \text{ MeV}}$$

Hence this reaction is exoergic ( $Q$  positive) and therefore the reaction does not require a threshold energy for this reaction to occur (same for all (n,  $\gamma$ ) and some (n, fission))